



Baxter Q-operator and functional relations



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ARTICLE INFO

Article history:

Received 30 December 2014

Accepted 2 February 2015

Available online 4 February 2015

Editor: L. Alvarez-Gaumé

ABSTRACT

We obtain the Baxter Q-operators in the $U_q(\widehat{sl}_2)$ invariant integrable models as a special limits of the quantum transfer matrices corresponding to different spins in the auxiliary space. We derive the Baxter equation from the well-known fusion relations for the transfer matrices. Our method is valid for an arbitrary integrable model corresponding to the quantum group $U_q(\widehat{sl}_2)$, for example for the XXZ-spin chain.

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1. Introduction

At present time the quantum groups and the universal R-matrix [1,2] are the powerful tools to study various aspects of quantum integrable models such as the Yang–Baxter equation, Baxter Q-operator and functional relations. Using the special q-oscillator representations of the quantum group $U_q(\widehat{sl}_2)$ for the auxiliary space the Baxter Q-operator [3] was constructed for the Conformal Field Theory [4] and the six-vertex model [5,6]. Using the explicit form of the universal R-matrix [7] the L-operator for the Baxter Q-operator was derived in [8]. The functional relations in the universal form for the $U_q(\widehat{sl}_2)$ -invariant integrable models were obtained in [9]. In the present paper we continue to study the Baxter Q-operator in various $U_q(\widehat{sl}_2)$ -invariant integrable models. We discuss the connection between the $j \rightarrow \pm\infty$ limits of the quantum transfer matrix Z_j and the Baxter Q-operators. First, we derive the correspondence between the limits of the spin- j representations of the quantum group and the q-oscillator representations. Second, we find the limits of the transfer matrices and show that it is given by the Q-operators. Finally, we obtain the Baxter equation as a limit of the functional relations in the universal form (fusion relations). Thus we show that the Q-operator could be derived without the knowledge of the q-oscillator representations.

2. Quantum group $U_q(\widehat{sl}_2)$

We begin with the definition of the algebra $U_q(\widehat{sl}_2)$ and the corresponding universal R-matrix. It is the complex associative algebra with the generators corresponding to the simple roots $e_i, f_i, h_i, i = 0, 1$, obeying the following defining relations:

$$[h_i, h_j] = 0, \quad [h_i, e_j] = a_{ij}e_j, \quad [h_i, f_j] = -a_{ij}f_j,$$

$$[e_i, f_j] = \delta_{ij} \frac{q^{h_i} - q^{-h_i}}{q - q^{-1}},$$

where a_{ij} is the Cartan matrix A which for the case of the algebra \widehat{sl}_2 equals

$$A = \begin{pmatrix} 2 & -2 \\ -2 & 2 \end{pmatrix}.$$

This definition should be supplemented by the Serre relations satisfied for all distinct i and j :

$$e_i^3 e_j - [3]_q e_i^2 e_j e_i + [3]_q e_i e_j e_i^2 - e_j e_i^3 = 0,$$

$$f_i^3 f_j - [3]_q f_i^2 f_j f_i + [3]_q f_i f_j f_i^2 - f_j f_i^3 = 0.$$

We do not need the grading element d in the present paper. The generators e_i, f_i, h_i correspond the simple roots vectors α, β according to the rule $e_0 = e_\beta, e_1 = e_\alpha$ and the same for f_i, h_i . Define the root vector $\delta = \alpha + \beta$. Then the root lattice consists of the root vectors $\alpha + m\delta, \beta + m\delta$ and $m\delta$ where $m \geq 1$. The corresponding generators can be obtained from $e_{\alpha,\beta}, f_{\alpha,\beta}$ using the formulas given in Ref. [7]. The algebra $U_q(\widehat{sl}_2)$ is a Hopf algebra with the comultiplication Δ defined by the equations

$$\begin{aligned} \Delta(e_i) &= e_i \otimes q^{h_i} + 1 \otimes e_i, & \Delta(f_i) &= f_i \otimes 1 + q^{-h_i} \otimes f_i, \\ \Delta(h_i) &= h_i \otimes 1 + 1 \otimes h_i. \end{aligned} \quad (1)$$

Now we introduce the universal R matrix which is an element in $U_q(\widehat{sl}_2) \otimes U_q(\widehat{sl}_2)$. This matrix satisfies the following simple relation:

$$\Delta'(g)R = R\Delta(g), \quad g \in U_q(\widehat{sl}_2), \quad (2)$$

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where $\Delta'(g)$ is obtained from $\Delta(g)$ by the permutation of the corresponding operators. The Hopf algebra is called quasitriangular if the following relations are satisfied:

$$(\Delta \otimes 1)R = R_{13}R_{23}, \quad (1 \otimes \Delta)R = R_{13}R_{12}, \quad (3)$$

where the notations $R_{12} = R \otimes 1$, $R_{23} = 1 \otimes R$ are used. From Eqs. (2), (3) it follows that the universal R -matrix satisfies the Yang–Baxter equation:

$$R_{12}R_{13}R_{23} = R_{23}R_{13}R_{12}. \quad (4)$$

The explicit expression for the universal R -matrix for the case of the algebra $U_q(\widehat{sl}_2)$ is given in [7]. For us it is sufficient to establish the following facts about it. The universal R -matrix can be represented in the form $R = \sum_i A_i \otimes B_i \in U_q(b_+) \otimes U_q(b_-)$, where b_{\pm} are the positive and the negative Borel subalgebras of the algebra \widehat{sl}_2 . More precisely it has the form $R = K \sum_i A_i \otimes B_i$, where $K = q^{(h_{\alpha} \otimes h_{\alpha})/2}$ and the operators A_i and B_i are some polynomials in the generators $e_{\alpha, \beta}$ and $f_{\alpha, \beta}$ respectively. Once the representations of the quantum group V_1, V_2 are specified the corresponding matrix R_{12} which is the building block in the framework of the Quantum Inverse Scattering Method is given by the projection of the universal R -matrix: $R_{12} = (V_1 \otimes V_2)R$.

3. Representations of quantum group and their correspondence

Let us introduce the representations of the quantum group which are required for the construction of the transfer-matrices and Q -operators according to the Quantum Inverse Scattering Method. First consider the representation $V_j(\lambda)$ which is obtained as a combination of Jimbo evaluation homomorphism and the standard finite-dimensional highest weight spin j representation of the algebra $U_q(\widehat{sl}_2)$. In fact consider the mapping

$$\begin{aligned} e_{\alpha} &= f, & f_{\alpha} &= e, & h_{\alpha} &= -h; \\ e_{\beta} &= \lambda e, & f_{\beta} &= \lambda^{-1} f, & h_{\beta} &= h; \end{aligned} \quad (5)$$

where e, f, h are the standard generators of $U_q(\widehat{sl}_2)$:

$$[h, e] = 2e, \quad [h, f] = -2f, \quad [e, f] = \frac{q^h - q^{-h}}{q - q^{-1}}.$$

The highest weight spin j representation is realized in the basis $|k\rangle$, $k = 0, 1, \dots, 2j$ in the following way:

$$\begin{aligned} f|k\rangle &= |k+1\rangle, & e|k\rangle &= [k]_q[2j+1-k]_q|k-1\rangle, \\ h|k\rangle &= (2j-2k)|k\rangle, \end{aligned} \quad (6)$$

where $[x]_q = (q^x - q^{-x})/(q - q^{-1})$. These formulas give the representation $V_j(\lambda)$. We denote the corresponding transfer-matrix by $Z_j(\lambda)$. We also need the representation $V_j^+(\lambda)$, where j is an arbitrary complex number, given by the same Eqs. (5), (6) except that the basis $|k\rangle$ is now infinite $k = 0, 1, \dots, \infty$. We denote the corresponding transfer matrix by $Z_j^+(\lambda)$.

To construct the Q -operators we will also need the two representations V^{\pm} of the positive Borel subalgebra $U_q(b_+)$ i.e. the subalgebra given by the positive generators $e_{\alpha, \beta}, h_{\alpha, \beta}$ and their descendants. These representations cannot be extended to the representations of the full quantum group. The representations V^{\pm} are build up from the so-called q -oscillator algebras [4]. We define two different operator algebras. The first one consists of the operators a^+, a, H_a obeying the following commutational relations:

$$\begin{aligned} [H_a; a^+] &= 2a^+, & [H_a; a] &= -2a, \\ qa^+a - \frac{1}{q}aa^+ &= \frac{\lambda}{(q - q^{-1})}. \end{aligned}$$

The second algebra consists of the operators b^+, b, H_b obeying the following commutational relations:

$$\begin{aligned} [H_b; b^+] &= 2b^+, & [H_b; b] &= -2b, \\ qbb^+ - \frac{1}{q}b^+b &= \frac{\lambda}{(q - q^{-1})}. \end{aligned}$$

One can choose the basis $|k\rangle$, $k = 0, 1, \dots, \infty$, such that the action of the operators is:

$$\begin{aligned} a^+|k\rangle &= |k+1\rangle, & a|k\rangle &= \frac{\lambda}{(q - q^{-1})^2} (1 - q^{2k}) |k-1\rangle, \\ H_a|k\rangle &= 2k|k\rangle \end{aligned} \quad (7)$$

for the first case and

$$\begin{aligned} b^+|k\rangle &= |k+1\rangle, & b|k\rangle &= \frac{\lambda}{(q - q^{-1})^2} (1 - q^{-2k}) |k-1\rangle, \\ H_b|k\rangle &= 2k|k\rangle \end{aligned} \quad (8)$$

for the second case. Then the q -oscillator representations V^{\pm} of the subalgebra $U_q(b_+)$ are defined by the following relations:

$$e_{\alpha} = a^+, \quad e_{\beta} = a, \quad h_{\alpha} = H_a, \quad h_{\beta} = -H_a \quad (9)$$

for the representation V^+ and

$$e_{\alpha} = b^+, \quad e_{\beta} = b, \quad h_{\alpha} = H_b, \quad h_{\beta} = -H_b \quad (10)$$

for the representation V^- . One can show that the operators corresponding to the auxiliary spaces given by the representations V^{\pm} are the Baxter Q -operators $Q^{\pm}(\lambda)$.

Now from Eqs. (5)–(10) one can see that the representations V^{\pm} can be obtained as a special limits of the representations $V_j(\lambda)$, $V_j^+(\lambda)$ (for example, see [9]). In fact one has the equations:

$$\begin{aligned} \lim_{j \rightarrow \infty} \lambda q^{-2j-1} [k]_q [2j+1-k]_q & \\ \rightarrow \frac{\lambda}{(q - q^{-1})^2} (1 - q^{-2k}), & \quad (V^-), \end{aligned} \quad (11)$$

$$\begin{aligned} \lim_{j \rightarrow -\infty} \lambda q^{2j+1} [k]_q [2j+1-k]_q & \\ \rightarrow \frac{\lambda}{(q - q^{-1})^2} (1 - q^{2k}), & \quad (V^+), \end{aligned} \quad (12)$$

which show that the representations V^{\pm} of the positive Borel subalgebra $U_q(b_+)$ can be obtained as a limits $j \rightarrow \pm\infty$ of the representation V_j^+ . Note that this correspondence cannot be extended to the representations of the full quantum group $U_q(\widehat{sl}_2)$. On the level of generators of quantum group the last equations imply the following substitution of generators:

$$e_{\alpha} = f \rightarrow a^+(\lambda), \quad e_{\beta} = \lambda q^{-2j-1} e(j) \rightarrow a(\lambda), \quad (V^+),$$

and

$$e_{\alpha} = f \rightarrow b^+(\lambda), \quad e_{\beta} = \lambda q^{2j+1} e(j) \rightarrow b(\lambda), \quad (V^-),$$

which will be used in the next section.

4. Tensor product of the representations and functional relations

Let us derive the functional relations in the universal form the representations of the quantum group. The simplest way to derive the Baxter TQ -equations and the so-called fusion relations for the product of two transfer matrices is to consider the tensor product of two representations V^{\pm} . For example, consider the tensor product $V^- \otimes V^+$ in detail. One should consider the product of the

two operators $Q^-(\lambda'')Q^+(\lambda')$. For each of the operators one has the following formal expression:

$$Q^\pm(\lambda) = \text{Tr}_{V^- \otimes V^+} R = \sum_i \sum_{k=0}^{\infty} q^{kS} \langle k | A_i(V^\pm) | k \rangle B_i, \quad (13)$$

where B_i are the operators which act in the quantum space and we denote by S the value of the generator $-h_\alpha$ in the quantum space ($S = S^z$ for the XXZ-spin chain). To evaluate the product of two Q -operators one can use the first of the relations (3) to get the following expression:

$$\begin{aligned} Q^-(\lambda'')Q^+(\lambda') \\ = \sum_i \sum_{n,l=0}^{\infty} q^{(n+l)S} \langle n, l | \Delta(A_i)(V^- \otimes V^+) | n, l \rangle B_i, \end{aligned} \quad (14)$$

where $|n, l\rangle = |n\rangle|l\rangle$ is the basis in the space $V^- \otimes V^+$ (the operators a^+, a act in the first space, and the operators b^+, b – in the second, see Section 3). Now one should choose the convenient basis in the space $|n, l\rangle$. Using Eqs. (1) one finds the following expressions:

$$\Delta(e_\alpha) = a^+ + b^+ q^{H_\alpha}, \quad \Delta(e_\beta) = a + b q^{-H_\alpha}.$$

The convenient basis is given by the equation:

$$|\phi_{k,m}\rangle = (\Delta(e_\alpha))^k |m\rangle|0\rangle.$$

Thus Eq. (14) can be represented in the form:

$$\begin{aligned} Q^-(\lambda'')Q^+(\lambda') \\ = \sum_i \sum_{k,m=0}^{\infty} q^{(k+m)S} \langle \phi_{k,m} | \Delta(A_i)(V^- \otimes V^+) | \phi_{k,m} \rangle B_i. \end{aligned} \quad (15)$$

The action of the operators $\Delta(e_{\alpha,\beta})$ on the states $|\phi_{k,m}\rangle$ has the following form:

$$\begin{aligned} \Delta(e_\alpha)|\phi_{k,m}\rangle &= |\phi_{k+1,m}\rangle, \\ \Delta(e_\beta)|\phi_{k,m}\rangle &= \lambda[k]_q [2j+1-k]_q |\phi_{k-1,m}\rangle + c_{k,m} |\phi_{k,m-1}\rangle \end{aligned} \quad (16)$$

for the spectral parameters $\lambda' = \lambda q^{-2j-1}$, $\lambda'' = \lambda q^{2j+1}$. The explicit form of the constants $c_{k,m}$ is not important. The commutativity of $Z_j(\lambda)$ and $Q^\pm(\lambda)$ follows from the universal Yang–Baxter equation (4), while the commutativity of $Q^\pm(\lambda)$ follows from the fact that these operators can be obtained as a special limits of the transfer matrices (see Section 5). From Eqs. (15), (16) we come to the transfer matrix $Z_j^+(\lambda)$ which has the following formal expression:

$$Z_j^+(\lambda) = \sum_i \sum_{k=0}^{\infty} q^{(k-j)S} \langle k | A_i(V_j^+) | k \rangle B_i. \quad (17)$$

Comparing this equation with Eq. (15) we obtain the following basic functional relation:

$$Z_j^+(\lambda) = \frac{1-q^S}{q^{jS}} Q^-(\lambda q^{2j+1}) Q^+(\lambda q^{-2j-1}). \quad (18)$$

It is convenient to redefine the operators $Q^\pm(\lambda)$ according to the equations:

$$Q^-(\lambda) = \lambda^{S/4} Q'^-(\lambda), \quad Q^+(\lambda) = \lambda^{-S/4} Q'^+(\lambda).$$

Then the relation (18) takes the form:

$$\begin{aligned} Z_j^+(\lambda) &= C Q'^-(\lambda q^{2j+1}) Q'^+(\lambda q^{-2j-1}), \\ C &= q^{S/2} (1 - q^S). \end{aligned} \quad (19)$$

The next step is to find the relation connecting the transfer-matrices $Z_j(\lambda)$ and $Z_j^+(\lambda)$. Taking into account Eqs. (6), (17) one can see that the sum over k from $2j+1$ to ∞ in Eq. (17) is equal to $Z_{-j-1}^+(\lambda)$. Thus we obtain:

$$\begin{aligned} Z_j^+(\lambda) &= Z_j(\lambda) + Z_{-j-1}^+(\lambda), \\ Z_j(\lambda) &= Z_j^+(\lambda) - Z_{-j-1}^+(\lambda). \end{aligned} \quad (20)$$

From these equations we obtain the following fundamental functional relation:

$$\begin{aligned} Z_j(\lambda) &= C \left(Q'^-(\lambda q^{2j+1}) Q'^+(\lambda q^{-2j-1}) \right. \\ &\quad \left. - Q'^-(\lambda q^{-2j-1}) Q'^+(\lambda q^{2j+1}) \right). \end{aligned} \quad (21)$$

From Eq. (21) one can derive all the other functional relations in the universal form. For example, taking into account the $Z_0(\lambda) = 1$, $Z_{1/2}(\lambda) = Z(\lambda)$ one can easily obtain the Baxter equation in the universal form:

$$Z(\lambda) Q'^\pm(\lambda) = Q'^\pm(\lambda q^2) + Q'^\pm(\lambda q^{-2}). \quad (22)$$

It is easy to obtain these relations considering the tensor product of the representations $V_{1/2} \otimes V^\pm$. In a similar way using (21), (22) one can obtain the following fusion relations for the transfer matrices:

$$Z(\lambda) Z_j(\lambda q^{-2j-1}) = Z_{j-1/2}(\lambda q^{-2j-2}) + Z_{j+1/2}(\lambda q^{-2j}) \quad (23)$$

and

$$Z(\lambda) Z_j(\lambda q^{2j+1}) = Z_{j-1/2}(\lambda q^{2j+2}) + Z_{j+1/2}(\lambda q^{2j}). \quad (24)$$

All the relations presented above are written down in the universal form i.e. for an arbitrary integrable model in the quantum space $U_q(b_-)$. The normalization of the operators corresponds to the appropriate projection of the universal R -matrix. In order to write down these operators in the polynomial form one should calculate the corresponding expressions for the universal R -matrix [8].

5. Limits of quantum transfer matrices as Baxter Q -operators

In this section we derive the limits $j \rightarrow \pm\infty$ of the transfer matrices $Z_j(\lambda)$ and $Z_j^+(\lambda)$ and show their relations to the Baxter Q -operators. Consider the transfer matrix $Z_j^+(\lambda)$. First, one can take the limits $j \rightarrow \pm\infty$ in Eq. (18) and take into account that $Q^\pm(0) = (1 - q^S)^{-1}$ (in order to avoid singularities at $S = 0$ one should introduce the twist angle ϕ , we will put $\phi = 0$ everywhere for simplicity and consider the case $\phi = 0$ in all the final expressions). Second, the same result can be obtained from Eqs. (11), (12). In fact the sum over k in Eq. (17) is determined by the low values of k (the sum in Eq. (13) converges). Thus both methods leads to the same result:

$$\lim_{j \rightarrow \infty} \left(Z_j^+(\lambda q^{-2j-1}) q^{jS} \right) = Q^-(\lambda), \quad (25)$$

$$\lim_{j \rightarrow -\infty} \left(Z_j^+(\lambda q^{2j+1}) q^{jS} \right) = Q^+(\lambda). \quad (26)$$

One can argue that in the case $j \rightarrow \infty$ the number k can take the values $k \sim 2j$ and Eq. (11) breaks down. However, we see from Eq. (18) that the contributions from the regions below the value $k = 2j+1$ and above this value cancel each other. Now let us consider the limits of the usual transfer matrix $Z_j(\lambda)$. It is convenient to represent Eq. (21) in the following form:

$$Z_j(\lambda) = \frac{1-q^S}{q^{jS}} Q^-(\lambda q^{2j+1}) Q^+(\lambda q^{-2j-1}) - \frac{1-q^S}{q^{(-j-1)S}} Q^-(\lambda q^{-2j-1}) Q^+(\lambda q^{2j+1}).$$

From this equation it is easy to obtain the following limits:

$$\begin{aligned} \lim_{j \rightarrow \infty} (Z_j(\lambda q^{-2j-1})) &= Q^-(\lambda) - Q^+(\lambda), \quad S = 0, \\ \lim_{j \rightarrow \infty} (Z_j(\lambda q^{-2j-1}) q^{(-j-1)S}) &= -Q^+(\lambda), \quad S > 0, \\ \lim_{j \rightarrow \infty} (Z_j(\lambda q^{-2j-1}) q^{jS}) &= Q^-(\lambda), \quad S < 0, \end{aligned} \quad (27)$$

and

$$\begin{aligned} \lim_{j \rightarrow -\infty} (Z_j(\lambda q^{2j+1})) &= Q^+(\lambda) - Q^-(\lambda), \quad S = 0, \\ \lim_{j \rightarrow -\infty} (Z_j(\lambda q^{2j+1}) q^{jS}) &= Q^+(\lambda), \quad S > 0, \\ \lim_{j \rightarrow -\infty} (Z_j(\lambda q^{2j+1}) q^{(-j-1)S}) &= -Q^-(\lambda), \quad S < 0. \end{aligned} \quad (28)$$

Note that the transfer matrix $Z_j(\lambda)$ at $j < 0$ is obtained as an analytical continuation of this function at $j = 0, 1/2, 1, \dots$. Consider the limits $j \rightarrow \infty$ in more details. First, consider the case $S = 0$. Clearly the contribution $Q^-(\lambda)$ comes from the low values of $k \sim 1$ in the sum over k in the formal equation of the type (17):

$$Z_j(\lambda) = \sum_i \sum_{k=0}^{2j} q^{(k-j)S} \langle k | A_i(V_j) | k \rangle B_i. \quad (29)$$

The term $-Q^+(\lambda)$ comes from the values $k \sim 2j$ in the sum over k in this equation. This fact can be confirmed by the following arguments. Consider Eq. (25). The sum over k is infinite for Z_j^+ . The contribution of the region $k \geq 2j+1$ in Eq. (17) equals $Q^+(\lambda)$. In fact, substituting $k = 2j+1+k'$, $k' \geq 0$ in Eq. (6) and taking the limit $j \rightarrow \infty$ we obtain exactly the operator $Q^+(\lambda)$. Taking into account Eqs. (17), (25), (29) one can see that the contribution of the region $k = 2j+1-k'$, $k' > 0$, is equal to $-Q^+(\lambda)$ in agreement with Eq. (27). One can come to the same conclusion by noticing that the contribution of the region above $k = 2j+1$ is equal to the limit of $Z_{-j-1}^+(\lambda q^{-2j-1})$ which can be easily evaluated. Thus the appearance of the term $-Q^+(\lambda)$ in Eq. (27) is explained. One can also calculate this contribution explicitly. In fact, taking $k = 2j+1-k'$, $k' \geq 0$, we see that the representation of the quantum group corresponding to this contribution is:

$$\begin{aligned} e_\alpha |k'\rangle &= |k'-1\rangle, \\ e_\beta |k'\rangle &= \frac{\lambda}{(q-q^{-1})^2} (1-q^{-2k'}) |k'+1\rangle, \end{aligned} \quad (30)$$

where $k' \geq 0$. From Eq. (30) one can see that the contribution of the i -th term in the decomposition $R = K \sum_i A_i \otimes B_i$ has the form

$$\sum_{N=0}^{N_i} C_N \sum_{k'=1}^{\infty} q^{-(S+2N)k'},$$

where C_N are some constants and N_i are the number of the generators e_β in A_i , while the contribution of the same term for the operator $Q^+(\lambda)$ equals

$$\sum_{N=0}^{N_i} C_N \sum_{k=0}^{\infty} q^{(S+2N)k}.$$

Comparing these two contributions one can easily see that the desired contribution is in fact equal to $-Q^+(\lambda)$. Next one can see that the limits (27) for $S > 0$ and $S < 0$ agree with the general formula (29). The case $S < 0$ is obvious while the limit in the case $S > 0$ follows from the arguments presented above. As for the limits (28) they are obviously related to the limits (27) via the relation $Z_j(\lambda) = -Z_{-j-1}(\lambda)$. Finally one can also see that the limits (27), (28) are in agreement with the relation (20). While Eqs. (25), (26) where written down in [10], Eqs. (27), (28) – are the main result of the present paper. We have found the limits of the transfer matrices $Z_j(\lambda)$, $Z_j^+(\lambda)$ both by the explicit calculations and from the functional relations (19), (21).

Now we show how to obtain the Baxter equation from the fusion relations without the knowledge of the explicit form of the Q -operators and their relation to the q -oscillator algebras. The point is that the fusion relations for the transfer matrices were obtained without the use of the quantum groups so that the Baxter Q -operator could be obtained without the knowledge of the q -oscillator representations. For example, consider the fusion relation (23). This equation can be represented in the following form:

$$\begin{aligned} Z(\lambda) Z_j(\lambda q^{-2j-1}) &= Z_{j-1/2}(\lambda q^{-2(j-1/2)-1} q^{-2}) + Z_{j+1/2}(\lambda q^{-2(j+1/2)-1} q^2). \end{aligned}$$

Considering the case $S < 0$, multiplying both sides of this equation by the factor q^{jS} and taking the limit $j \rightarrow \infty$ we obtain the equation:

$$Z(\lambda) Q^-(\lambda) = Q^-(\lambda q^{-2}) q^{S/2} + Q^-(\lambda q^2) q^{-S/2}.$$

It is easy to prove that in all the other cases we get the same equation. In the same way from the fusion relation (24) we obtain

$$\begin{aligned} Z(\lambda) Z_j(\lambda q^{2j+1}) &= Z_{j-1/2}(\lambda q^{2(j-1/2)+1} q^2) + Z_{j+1/2}(\lambda q^{2(j+1/2)+1} q^{-2}), \end{aligned}$$

which in the limit $j \rightarrow -\infty$ leads to the equation:

$$Z(\lambda) Q^+(\lambda) = Q^+(\lambda q^{-2}) q^{-S/2} + Q^+(\lambda q^2) q^{S/2}.$$

These equations are equivalent to the usual equations for the standard Baxter Q -operators $Q'^{\pm}(\lambda)$:

$$Z(\lambda) Q'^{\pm}(\lambda) = Q'^{\pm}(\lambda q^2) + Q'^{\pm}(\lambda q^{-2}).$$

Thus we obtained the Baxter Q -operators without the knowledge about their connection with the q -oscillator algebras. Surely this is connected with the relations of the representations of quantum group given in Section 3.

6. Conclusion

In conclusion, we obtained the Baxter Q -operators in the quantum group $U_q(\mathfrak{sl}_2)$ invariant integrable models as a limits of the quantum transfer matrices corresponding to different spins in the auxiliary space. We showed that the Baxter equation and the form of the Q -operator follow from the usual fusion functional relations for the quantum transfer matrices. All the relations are represented in the universal form and valid for an arbitrary quantum integrable model in the quantum space, for example, for the XXZ-spin chain. In order to find the functional relations for the concrete model, one should calculate the universal R -matrix in a given representa-

tion [9], or to find the R - and L -operators as a polynomials in the spectral parameter.

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